

Глава 2. Некоторые классические задачи математической физики

§5. Уравнение $\Delta u + cu = -f$ в неограниченной области

Парциальные условия излучения ($e^{-i\omega t}$).

$$\int_0^b \left\{ \frac{\partial u}{\partial x} + i\gamma_n^{(1)} u \right\}_{x=0} \psi_n(y) dy = 2i\gamma_{n_0} A_{n_0} \delta_{nn_0} \quad (44) \quad (n = 1, 2, \dots)$$

Падающая(правая) волна

$$u_{n_0}(x, y) = A_{n_0} e^{i\gamma_{n_0} x} \psi_{n_0}(y) \Rightarrow$$

$$\begin{aligned} \int_0^b \left\{ \frac{\partial u_{n_0}}{\partial x} + i\gamma_n^{(1)} u_{n_0} \right\}_{x=0} \psi_n(y) dy &= \int_0^b \left\{ i\gamma_{n_0}^{(1)} A_{n_0} + i\gamma_n^{(1)} A_{n_0} \right\} \psi_{n_0}(y) \psi_n(y) dy = \\ &= \left\{ i\gamma_{n_0}^{(1)} A_{n_0} + i\gamma_n^{(1)} A_{n_0} \right\} \boxed{\int_0^b \psi_{n_0}(y) \psi_n(y) dy} = \\ &= \delta_{nn_0} \end{aligned}$$

$$= 2i\gamma_{n_0} A_{n_0} \delta_{nn_0} \Rightarrow \text{падающая волна } u_{n_0}(x, y) \text{ проходит сечение } x=0$$

Любая правая волна при $n \neq n_0$ пройти сечение $x=0$ не может:

$$\begin{aligned} (\rightarrow) \\ u_k(x, y) = A_k e^{i\gamma_k x} \psi_k(y), \quad k \neq n_0 \Rightarrow \end{aligned}$$

$$\begin{aligned} \int_0^b \left\{ \frac{\partial u_k}{\partial x} + i\gamma_n^{(1)} u_k \right\}_{x=0} \psi_n(y) dy &= \left\{ i\gamma_k^{(1)} A_k + i\gamma_n^{(1)} A_k \right\}_{x=0} \boxed{\int_0^b \psi_n(y) \psi_k(y) dy} = \\ &= \delta_{nk} \\ &= 2i\gamma_k^{(1)} A_k = 2i\gamma_{n_0}^{(1)} A_{n_0} \boxed{\delta_{kn_0}} = 0 \Rightarrow A_k = 0, k \neq n_0 \end{aligned}$$

$$\begin{aligned} (\leftarrow) \\ \text{Любая левая волна проходит сечение } x=0: u_k(x, y) = B_k e^{-i\gamma_k x} \psi_k(y) \Rightarrow \end{aligned}$$

$$\int_0^b \left\{ \frac{\partial u_k}{\partial x} + i\gamma_n^{(1)} u_k \right\}_{x=0} \psi_n(y) dy =$$

$$= \left\{ -i\gamma_k^{(1)} B_k + i\gamma_n^{(1)} B_n \right\}_{x=0} \boxed{\int_0^b \psi_n(y) \psi_k(y) dy} =$$

$$= (-i\gamma_n^{(1)} B_n + i\gamma_n^{(1)} B_n) = 0$$

Квадрупольный излучатель

$$p(r, \theta) = \sum_{n=0}^{\infty} C_n \zeta_n^{(1)}(kr) P_n(\cos \theta) \quad (61)$$

$$\zeta_n^{(1)}(kr) = \sqrt{\frac{\pi}{2kr}} H_{n+\frac{1}{2}}^{(1)}(kr) \quad (62)$$

$$\left. \frac{\partial p}{\partial r} \right|_{r=a} = i\omega \rho_0 v_0 P_2(\cos \theta) \quad (58) \Rightarrow$$

$$\sum_{n=0}^{\infty} C_n k \zeta_n^{(1)'}(ka) P_n(\cos \theta) = i\omega \rho_0 v_0 P_2(\cos \theta)$$

$$C_2 k \zeta_2^{(1)'}(ka) = i\omega \rho_0 v_0 \Rightarrow$$

$$C_2 = i \left(\frac{\omega}{k} \right)^c \frac{\rho_0 v_0}{\zeta_2^{(1)'}(ka)} = ic \rho_0 v_0 \frac{1}{\zeta_2^{(1)'}(ka)} \Rightarrow$$

$$p(r, \theta) = ic \rho_0 v_0 \frac{\zeta_2^{(1)}(kr)}{\zeta_2^{(1)'}(ka)} P_2(\cos \theta) \quad (64)$$

$$v_r = \frac{1}{i\omega \rho_0} \frac{\partial p}{\partial r} \quad (55) \Rightarrow$$

= \omega

$$v_r(r, \theta) = \frac{1}{i\omega \rho_0} \overbrace{ic \rho_0 v_0 k} \frac{\zeta_2^{(1)'}(kr)}{\zeta_2^{(1)'}(ka)} P_2(\cos \theta) =$$

$$= v_0 \frac{\zeta_2^{(1)'}(kr)}{\zeta_2^{(1)'}(ka)} P_2(\cos \theta) \quad (65)$$

$$ka = \frac{2\pi}{\lambda} a \ll 1, \quad kr = \frac{2\pi}{\lambda} r \gg 1$$

$$\begin{aligned} p(r, \theta) &= ic\rho_0 v_0 \frac{\zeta_2^{(1)'}(kr)}{\zeta_2^{(1)'}(ka)} P_2(\cos \theta) = ic\rho_0 v_0 \frac{-\frac{1}{kr} e^{ikr} e^{-i\frac{\pi}{2}}}{i \frac{9}{(ka)^4}} P_2(\cos \theta) = \\ &= \frac{ic\rho_0 v_0 k^4 a^4}{i9kr} e^{ikr} e^{-i\frac{\pi}{2}} P_2(\cos \theta) = -\frac{c\rho_0 v_0 k^3 a^4}{9r} e^{ikr} e^{-i\frac{\pi}{2}} P_2(\cos \theta) \quad (67) \end{aligned}$$

$$\begin{aligned} v_r(r, \theta) &= v_0 \frac{\zeta_2^{(1)'}(kr)}{\zeta_2^{(1)'}(ka)} P_2(\cos \theta) = v_0 \frac{-\frac{i}{kr} e^{ikr} e^{-i\frac{\pi}{2}}}{i \frac{9}{(ka)^4}} P_2(\cos \theta) = \\ &= -\frac{v_0 k^4 a^4}{9kr} e^{ikr} e^{-i\frac{\pi}{2}} P_2(\cos \theta) = -\frac{v_0 k^3 a^4}{9r} e^{ikr} e^{-i\frac{\pi}{2}} P_2(\cos \theta) \quad (68) \end{aligned}$$

$$\begin{aligned} \bar{P}(r, \theta, t) &= \text{Re}(p(r, \theta) e^{-i\omega t}) = \text{Re} \left(-\frac{c\rho_0 v_0 k^3 a^4}{9r} P_2(\cos \theta) e^{-i\left(\omega t - kr + \frac{\pi}{2}\right)} \right) = \\ &= -\frac{c\rho_0 v_0 k^3 a^4}{9r} P_2(\cos \theta) \cos \left(\omega t - kr + \frac{\pi}{2} \right) \end{aligned}$$

$$\begin{aligned}\bar{V}_r(r, \theta, t) &= \text{Re}(v_r(r, \theta)e^{-i\omega t}) = \text{Re}\left(-\frac{v_0 k^3 a^4}{9r} P_2(\cos \theta) e^{-i\left(\omega t - kr + \frac{\pi}{2}\right)}\right) = \\ &= -\frac{v_0 k^3 a^4}{9r} P_2(\cos \theta) \cos\left(\omega t - kr + \frac{\pi}{2}\right)\end{aligned}$$

$$\begin{aligned}\bar{Y} &= \frac{1}{T} \int_0^T \bar{P}(r, \theta, t) \bar{V}_r(r, \theta, t) dt = \frac{c \rho_0 v_0^2 k^6 a^8 P_2^2(\cos \theta)}{81 r^2} \frac{1}{T} \int_0^T \cos^2\left(\omega t - kr + \frac{\pi}{2}\right) dt = \\ &= \frac{1}{T} \int_0^T \cos^2\left(\omega t - kr + \frac{\pi}{2}\right) dt = \frac{1}{2T} \int_0^T (1 + \cos 2\left(\omega t - kr + \frac{\pi}{2}\right)) dt = \frac{1}{2} \Rightarrow\end{aligned}$$

$$\bar{Y} = \frac{c \rho_0 v_0^2 k^6 a^8 P_2^2(\cos \theta)}{162 r^2}$$

$$\begin{aligned}\Pi &= r^2 \int_0^{2\pi} \int_0^\pi \bar{Y} \sin \theta d\theta d\varphi = \frac{2\pi c \rho_0 v_0^2 k^6 a^8}{162} \int_\pi^0 P_2^2(\cos \theta) d \cos \theta = \\ &= \frac{\pi c \rho_0 v_0^2 k^6 a^8}{81} \int_{-1}^1 P_2^2(x) dx = \frac{\pi c \rho_0 v_0^2 k^6 a^8}{81} \|P_2\|^2 = \\ &= \frac{\pi c \rho_0 v_0^2 k^6 a^8}{81} \frac{2}{2 \cdot 2 + 1} = \frac{2\pi c \rho_0 v_0^2 k^6 a^8}{405}\end{aligned}$$

\bar{Y} - интенсивность излучения квадруполь,

Π - мощность излучения квадруполь.